

On the resilience of scalar and vector vortex modes in turbulence

MITCHELL A. COX,^{1,2} CARMELO ROSALES-GUZMÁN,² MARTIN P. J. LAVERY,³ DANIEL J. VERSFELD¹ AND ANDREW FORBES^{2,*}

¹*School of Electrical and Information Engineering, University of the Witwatersrand, Johannesburg 2050, South Africa*

²*School of Physics, University of the Witwatersrand, Johannesburg 2050, South Africa*

³*School of Engineering, University of Glasgow, Lanarkshire, G12 8LT, UK*

**andrew.forbes@wits.ac.za*

Abstract: Free-space optical communication with spatial modes of light has become topical due to the possibility of dramatically increasing communication bandwidth via Mode Division Multiplexing (MDM). While both scalar and vector vortex modes have been used as transmission bases, it has been suggested that the latter is more robust in turbulence. Using orbital angular momentum as an example, we demonstrate theoretically and experimentally that the crosstalk due to turbulence is the same in the scalar and vector basis sets of such modes. This work brings new insights about the behaviour of vector and scalar modes in turbulence, but more importantly it demonstrates that when considering optimal modes for MDM, the choice should not necessarily be based on their vectorial nature.

© 2016 Optical Society of America

OCIS codes: (060.2605) Free-space optical communication; (010.1300) Atmospheric propagation; (260.5430) Polarization; (030.7060) Turbulence.

References and links

1. S. Berdagué and P. Facq, "Mode division multiplexing in optical fibers," *Appl. Opt.* **21**, 1950 (1982).
2. D. J. Richardson, "Filling the Light Pipe," *Science* **330**, 327 (2010).
3. A. E. Willner, H. Huang, Y. Yan, Y. Ren, N. Ahmed, G. Xie, C. Bao, L. Li, Y. Cao, Z. Zhao, J. Wang, M. P. J. Lavery, M. Tur, S. Ramachandran, A. F. Molisch, N. Ashrafi, and S. Ashrafi, "Optical communications using orbital angular momentum beams," *Adv. Opt. Photon.* **7**, 66 (2015).
4. A. Forbes, A. Dudley, and M. McLaren, "Creation and detection of optical modes with spatial light modulators," *Adv. Opt. Photon.* **8**, 200 (2016).
5. G. C. G. Berkhout, M. P. J. Lavery, J. Courtial, M. W. Beijersbergen, and M. J. Padgett, "Efficient sorting of orbital angular momentum states of light," *Phys. Rev. Lett.* **105**, 153601 (2010).
6. M. N. O'Sullivan, M. Mirhosseini, M. Malik, and R. W. Boyd, "Near-perfect sorting of orbital angular momentum and angular position states of light," *Opt. Express* **20**, 24444 (2012).
7. G. Gibson, J. Courtial, M. J. Padgett, M. Vasnetsov, V. Pas'ko, S. M. Barnett, and S. Franke-Arnold, "Free-space information transfer using light beams carrying orbital angular momentum," *Opt. Express* **12**, 5448 (2004).
8. M. Krenn, R. Fickler, M. Fink, J. Handsteiner, M. Malik, T. Scheidl, R. Ursin, and A. Zeilinger, "Communication with spatially modulated light through turbulent air across Vienna," *New J. Phys.* **16**, 113028 (2014).
9. J. Wang, J.-Y. Yang, I. M. Fazal, N. Ahmed, Y. Yan, H. Huang, Y. Ren, Y. Yue, S. Dolinar, M. Tur, and A. E. Willner, "Terabit free-space data transmission employing orbital angular momentum multiplexing," *Nat. Photon.* **6**, 488 (2012).
10. G. Funes, M. Vial, and J. A. Anguita, "Orbital-angular-momentum crosstalk and temporal fading in a terrestrial laser link using single-mode fiber coupling," *Opt. Express* **23**, 23133–42 (2015).
11. R. Brüning, B. Ndagano, and M. McLaren, "Data transmission with twisted light through a free-space to fiber optical communication link," *J. Opt.* **18**, 03LT01 (2016).
12. M. Andersson, E. Berglind, and G. Björk, "Orbital angular momentum modes do not increase the channel capacity in communication links," *New J. Phys.* **17**, 043040 (2015).
13. N. Zhao, X. Li, G. Li, and J. M. Kahn, "Capacity limits of spatially multiplexed free-space communication," *Nat. Photon.* **9**, 822 (2015).
14. J. A. Anguita, M. A. Neifeld, and B. V. Vasic, "Turbulence-induced channel crosstalk in an orbital angular momentum-multiplexed free-space optical link," *Appl. Opt.* **47**, 2414 (2008).
15. L. H.-C. Pu Ji-Xiong, Wang Tao and L. Cheng-Liang, "Propagation of cylindrical vector beams in a turbulent atmosphere," *Chin. Phys. B* **19**, 89201 (2010).
16. R. Chen, Y. Dong, F. Wang, and Y. Cai, "Statistical properties of a cylindrical vector partially coherent beam in turbulent atmosphere," *Appl. Phys. B* **112**, 247–259 (2013).

17. C. Chen, H. Yang, S. Tong, and Y. Lou, "Changes in orbital-angular-momentum modes of a propagated vortex Gaussian beam through weak-to-strong atmospheric turbulence," *Opt. Express* **24**, 6959 (2016).
18. B. Rodenburg, M. P. J. Lavery, M. Malik, M. N. O'Sullivan, M. Mirhosseini, D. J. Robertson, M. Padgett, and R. W. Boyd, "Influence of atmospheric turbulence on states of light carrying orbital angular momentum," *Opt. Lett.* **37**, 3735 (2012).
19. G. Milione, M. P. J. Lavery, H. Huang, Y. Ren, G. Xie, T. A. Nguyen, E. Karimi, L. Marrucci, D. A. Nolan, R. R. Alfano, and Others, "4× 20 Gbit/s mode division multiplexing over free space using vector modes and a q-plate mode (de) multiplexer," *Opt. Lett.* **40**, 1980 (2015).
20. Y. Ren, H. Huang, G. Xie, N. Ahmed, Y. Yan, B. I. Erkmen, N. Chandrasekaran, M. P. J. Lavery, N. K. Steinhoff, M. Tur, and Others, "Atmospheric turbulence effects on the performance of a free space optical link employing orbital angular momentum multiplexing," *Opt. Lett.* **38**, 4062 (2013).
21. Y. Ren, Z. Wang, P. Liao, L. Li, G. Xie, H. Huang, Z. Zhao, Y. Yan, N. Ahmed, A. Willner, M. P. J. Lavery, N. Ashrafi, S. Ashrafi, R. Bock, M. Tur, I. B. Djordjevic, M. A. Neifeld, and A. E. Willner, "Experimental characterization of a 400 Gbit/s orbital angular momentum multiplexed free-space optical link over 120 m," *Opt. Lett.* **41**, 622 (2016).
22. S. M. Zhao, J. Leach, L. Y. Gong, J. Ding, and B. Y. Zheng, "Aberration corrections for free-space optical communications in atmosphere turbulence using orbital angular momentum states," *Opt. Express* **20**, 452 (2012).
23. H. Huang, Y. Cao, G. Xie, Y. Ren, Y. Yan, C. Bao, N. Ahmed, M. A. Neifeld, S. J. Dolinar, and A. E. Willner, "Crosstalk mitigation in a free-space orbital angular momentum multiplexed communication link using 4x4 MIMO equalization," *Opt. Lett.* **39**, 4360 (2014).
24. Y. Ren, G. Xie, H. Huang, N. Ahmed, Y. Yan, L. Li, C. Bao, M. P. J. Lavery, M. Tur, M. A. Neifeld, R. W. Boyd, J. H. Shapiro, and A. E. Willner, "Adaptive-optics-based simultaneous pre- and post-turbulence compensation of multiple orbital-angular-momentum beams in a bidirectional free-space optical link," *Optica* **1**, 376 (2014).
25. G. Milione, H. I. Sztul, D. A. Nolan, and R. R. Alfano, "Higher-order poincaré sphere, stokes parameters, and the angular momentum of light," *Phys. Rev. Lett.* **107**, 53601 (2011).
26. G. Milione, T. A. Nguyen, J. Leach, D. A. Nolan, and R. R. Alfano, "Using the nonseparability of vector beams to encode information for optical communication," *Opt. Lett.* **40**, 4887 (2015).
27. Y. Gu, O. Korotkova, and G. Gbur, "Scintillation of nonuniformly polarized beams in atmospheric turbulence," *Opt. Lett.* **34**, 2261 (2009).
28. Y. Gu and G. Gbur, "Reduction of turbulence-induced scintillation by nonuniformly polarized beam arrays," *Opt. Lett.* **37**, 1553–1555 (2012).
29. W. Cheng, J. W. Haus, and Q. Zhan, "Propagation of vector vortex beams through a turbulent atmosphere," *Opt. Express* **17**, 17829 (2009).
30. L. Allen, M. W. Beijersbergen, R. J. C. Spreeuw, and J. P. Woerdman, "Orbital angular momentum of light and the transformation of Laguerre-Gaussian laser modes," *Phys. Rev. A* **45**, 8185 (1992).
31. D. Naidoo, F. S. Roux, A. Dudley, I. Litvin, B. Piccirillo, L. Marrucci, and A. Forbes, "Controlled generation of higher-order Poincaré sphere beams from a laser," *Nat. Photon.* **10**, 327 (2016).
32. M. Charnotskii, "Propagation of polarized waves in inhomogeneous media," *J. Opt. Soc. Am. A* **33**, 1385 (2016).
33. L. Marrucci, C. Manzo, and D. Paparo, "Optical spin-to-orbital angular momentum conversion in inhomogeneous anisotropic media," *Phys. Rev. Lett.* **96**, 163905 (2006).
34. R. G. Lane, A. Glindemann, and J. C. Dainty, "Simulations of a Kolmogorov phase screen," *Waves in Random Media* **2**, 209 (1992).
35. J. Liu and J. Wang, "Demonstration of polarization-insensitive spatial light modulation using a single polarization-sensitive spatial light modulator," *Sci. Rep.* **5**, 9959 (2015).
36. T. Kaiser, D. Flamm, S. Schröter, and M. Duparré, "Complete modal decomposition for optical fibers using CGH-based correlation filters," *Opt. Express* **17**, 9347 (2009).
37. V. P. Aksenov, V V Kolosov, G. A. Filimonov, and C. E. Pogutsa, "Orbital angular momentum of a laser beam in a turbulent medium: preservation of the average value and variance of fluctuations," *J. Opt.* **18**, 054013 (2016).

1. Introduction

Mode Division Multiplexing (MDM) was first suggested more than three decades ago as a technique to increase the bandwidth in fibre-based optical communications [1]. The realisation in recent years that current optical fibre communication systems will undergo a "capacity crunch" in the near future has renewed interest in the research [2]. Along this line, Orbital Angular Momentum (OAM) stands out as the mode of choice due to its topical nature and ease of measurement [3–6]. Successful demonstrations in both optical fibres and free space [3, 7–11] suggest the viability of OAM-based optical communication systems, even though its capacity limits still remain controversial [12, 13].

A limiting factor to the deployment of free-space optical communications is turbulent atmospheric conditions which lead to aberrations of the optical wave-front of a propagating beam.

These aberrations arise from localised changes in temperature and pressure, causing spatial variations in the refractive index of the atmosphere [14–17]. In MDM, turbulence causes degradation of the orthogonal spatial modes which results in the spreading of power into neighbouring modes. In the case of an MDM communication system, this so-called crosstalk imposes limits on the achievable channel capacity [18–21]. Investigation into the mitigation of the effects of atmospheric turbulence on beams carrying OAM have shown very promising results, where techniques employing specific modal selection, adaptive optics and digital signal processing have been demonstrated [22–24].

Vector vortex modes are non-separable states of light in which polarisation and OAM are coupled, which results in an inhomogeneous polarisation distribution [25]. Recently, the use of vector modes was proposed as an alternative to scalar modes to encode information [26]. It has been shown numerically that non-uniformly polarised beams such as a vector vortex beams are analogous to partially coherent beams, in as far as their resilience to atmospheric turbulence is concerned [27, 28]. Another study has postulated that the polarisation distribution of a vector vortex beam is maintained even after its intensity distribution has degraded and thus a portion of the information encoded in polarisation is still present [29]. From these studies it has been inferred that vector vortex beams are more resilient to turbulence as compared to their scalar counterparts [27–29].

Here we show that Cylindrical Vector Vortex (CVV) modes are not more resilient to atmospheric turbulence than their scalar (OAM) counterparts. We confirm this experimentally by measuring the crosstalk between basis elements of both mode sets, perturbed by Kolmogorov, thin phase screens encoded on a polarisation invariant Spatial Light Modulator (SLM). In comparing the modal crosstalk induced in each case, it is determined that although the crosstalk between modes within the scalar vortex and CVV bases is distributed differently, the total crosstalk is in fact identical within experimental error. The coupling between OAM and polarization in CVV beams is not sufficient to make their phase variation less susceptible to atmospheric turbulence when compared to circularly polarized scalar vortex beams.

2. Theory

A scalar basis set can be constructed by combining the degree of freedom provided by circular polarisation with the degrees of freedom given by the infinite set of OAM modes, illustrated in Fig. 1. These modes can be written, using Dirac notation for brevity, as:

$$|R^+\rangle = |\ell\rangle |R\rangle, \quad (1a) \quad |R^-\rangle = |-\ell\rangle |R\rangle, \quad (1c)$$

$$|L^+\rangle = |\ell\rangle |L\rangle, \quad (1b) \quad |L^-\rangle = |-\ell\rangle |L\rangle, \quad (1d)$$

where, R and L denotes right and left circular polarization, respectively, and $\ell \in \mathbb{Z}$ relates to the amount of orbital angular momentum, $\ell\hbar$, per photon [30]. The elements of this set are orthogonal to each other.

CVV modes are non-separable states of light where polarization and OAM modes are coupled. A basis set of four orthogonal CVV modes can be constructed as linear combinations of the scalar mode set in Eq. (1). The resulting modes carry no overall angular momentum and their transverse polarization state is not constant, illustrated in the bottom row of Fig. 1. Akin to the scalar case, these modes can be represented in the circular polarization basis using Dirac notation as:

$$|TM\rangle = 2^{-1/2}(|\ell\rangle |R\rangle + |-\ell\rangle |L\rangle), \quad (2a) \quad |HE^e\rangle = 2^{-1/2}(|\ell\rangle |L\rangle + |-\ell\rangle |R\rangle), \quad (2c)$$

$$|TE\rangle = 2^{-1/2}(|\ell\rangle |R\rangle - |-\ell\rangle |L\rangle), \quad (2b) \quad |HE^o\rangle = 2^{-1/2}(|\ell\rangle |L\rangle - |-\ell\rangle |R\rangle). \quad (2d)$$

Each mode in a given set (vector and scalar) is orthogonal with the others in the set, and

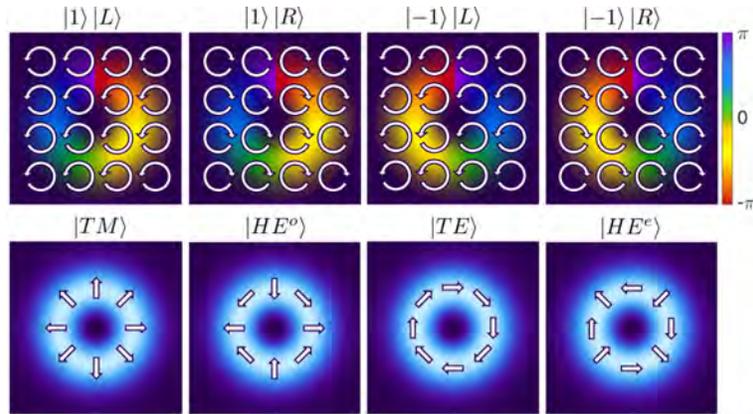


Fig. 1. Illustration of the scalar (top) and CVV (bottom) modes described in Eq. (1) and Eq. (2) respectively, with arrows indicating polarisation distribution which is constant for scalar and variable for CVV modes, for $\ell = \pm 1$. The phase of the four scalar modes in the transverse plane increases from $-\pi$ (red) to π (blue) for positive ℓ 's and decreases in the opposite direction for negative ℓ 's whereas the phase for CVV modes is a superposition of both.

mutually unbiased across the sets. Moreover, each of the modes described by Eq. (1) and (2) can be represented on a high order Poincaré sphere [25, 31].

2.1. Inter-mode crosstalk

In the presence of atmospheric turbulence, the polarisation of a beam is not affected because the atmosphere is not birefringent, which can readily be validated by tilting ones head (e.g., Kolmogorov turbulence assumes a homogeneous and isotropic atmosphere) [32]. The spatial degree of freedom, however, experiences aberrations which result in the coupling of modes into neighbouring modes, degrading their orthogonality. The amount of mode coupling or crosstalk is dependent on the strength of the turbulence in the channel, which can be expressed, in general, by:

$$|\ell\rangle \xrightarrow{\text{turb.}} \sum_{\ell'} p_{\ell-\ell'} |\ell'\rangle, \quad (3)$$

where $p_{\ell-\ell'}$ are the mode coupling weightings described by some distribution (e.g., p_0 would represent the modal power in the original OAM mode). This is used to find general expressions for what a given mode propagating through turbulence will transform into. For instance, the final state of the scalar mode $|R^+\rangle$ will be given by:

$$|R^+\rangle \xrightarrow{\text{turb.}} p_0 |\ell\rangle |R\rangle + p_{2\ell} |-\ell\rangle |R\rangle = |R_{\text{turb.}}^+\rangle \quad (4)$$

Analogous expressions can be found for $|R^-\rangle$, $|L^+\rangle$ and $|L^-\rangle$. Similarly, applying Equation 3 to a vector mode $|TM\rangle$:

$$|TM\rangle \xrightarrow{\text{turb.}} p_0 |\ell\rangle |R\rangle + p_0 |-\ell\rangle |L\rangle + p_{-2\ell} |\ell\rangle |L\rangle + p_{2\ell} |-\ell\rangle |R\rangle = |TM_{\text{turb.}}\rangle. \quad (5)$$

Analogous relations for $|TE\rangle$, $|HE^c\rangle$ and $|HE^o\rangle$ can also be found.

The crosstalk for each input mode is the magnitude of the inner product between each input mode and all the expected output modes after turbulence. By way of example, the inner product

of the scalar mode $|R^+\rangle$ with the mode $|R_{turb}^+\rangle$, will be:

$$\|\langle R^+ | R_{turb}^+ \rangle\|^2 = \|\langle \ell | \langle R | (p_0 |\ell\rangle |R\rangle + p_{2\ell} |-\ell\rangle |R\rangle) \|^2 = \|p_0\|^2. \quad (6)$$

Similarly, the inner product of the vector mode $|TM\rangle$ with the mode $|TM_{turb}\rangle$ will be:

$$\begin{aligned} \|\langle TM | TM_{turb} \rangle\|^2 = & \left\| \left(\frac{1}{\sqrt{2}} (\langle \ell | \langle R | + \langle -\ell | \langle L |) (p_0 |\ell\rangle |R\rangle + p_{2\ell} |-\ell\rangle |L\rangle + p_{-2\ell} |\ell\rangle |L\rangle \right. \right. \\ & \left. \left. + p_{2\ell} |-\ell\rangle |R\rangle) \right) \right\|^2 = \|p_0\|^2 \end{aligned} \quad (7)$$

The crosstalk terms for all of the modes can be summarised in the form of a matrix as:

$$M_{scalar} = \begin{pmatrix} \|p_0\|^2 & 0 & \|p_{-2\ell}\|^2 & 0 \\ 0 & \|p_0\|^2 & 0 & \|p_{-2\ell}\|^2 \\ \|p_{2\ell}\|^2 & 0 & \|p_0\|^2 & 0 \\ 0 & \|p_{2\ell}\|^2 & 0 & \|p_0\|^2 \end{pmatrix}, \quad (8)$$

$$M_{vector} = \begin{pmatrix} \|p_0\|^2 & \frac{\|p_{-2\ell} + p_{2\ell}\|^2}{4} & 0 & \frac{\|p_{-2\ell} - p_{2\ell}\|^2}{4} \\ \frac{\|p_{-2\ell} + p_{2\ell}\|^2}{4} & \|p_0\|^2 & \frac{\|p_{-2\ell} - p_{2\ell}\|^2}{4} & 0 \\ 0 & \frac{\|p_{-2\ell} - p_{2\ell}\|^2}{4} & \|p_0\|^2 & \frac{\|p_{-2\ell} + p_{2\ell}\|^2}{4} \\ \frac{\|p_{-2\ell} - p_{2\ell}\|^2}{4} & 0 & \frac{\|p_{-2\ell} + p_{2\ell}\|^2}{4} & \|p_0\|^2 \end{pmatrix}. \quad (9)$$

Notice that the terms in the diagonal are identical both cases. Under zero turbulence, where the mode coupling weightings for the crosstalk terms can be assumed to be zero ($p_{2\ell} = p_{-2\ell} = 0$), only the terms in the diagonal remain. This is consistent with the notion that when propagating in free-space with no aberrations induced by atmospheric turbulence, the output mode equals the input mode.

The total crosstalk, N , in each case can be computed by adding the off-diagonal elements as:

$$N_{vector} = \sum_{i \neq j} M_{vector} = (\|p_{2\ell} + p_{-2\ell}\|^2 + \|p_{2\ell} - p_{-2\ell}\|^2) = 2(\|p_{-2\ell}\|^2 + \|p_{2\ell}\|^2) \quad (10)$$

$$N_{scalar} = \sum_{i \neq j} M_{scalar} = 2(\|p_{-2\ell}\|^2 + \|p_{2\ell}\|^2) \quad (11)$$

which therefore leads to our claim that the total crosstalk in each case is in fact identical:

$$N_{scalar} = N_{vector} = 2(\|p_{-2\ell}\|^2 + \|p_{2\ell}\|^2) \quad (12)$$

3. Experiment

Figure 2 shows the experimental setup used to corroborate our theoretical findings. It consisted of three main stages: generation of scalar and CVV modes, turbulence using a Spatial Light Modulator (SLM) and finally, detection. In the first stage, both scalar and CVV modes were generated using a q -plate ($q = 1/2$) in conjunction with polarisers and wave plates [33]. The configuration of each element is detailed in the Appendix. For the second stage, a HoloEye Pluto SLM was used (PLUTO-VIS, 1920×1080 pixels with $8 \mu\text{m}$ pitch, calibrated for a 2π maximum phase shift at 633 nm) to simulate atmospheric turbulence. Random Kolmogorov phase screens were used, specified by their Strehl Ratio (SR) [34]. Since this SLM is only able to modulate horizontally polarised light, a polarisation invariant arrangement was implemented, as illustrated in Fig. 2, following the approach in [35]. In the last stage, the perturbed modes were detected

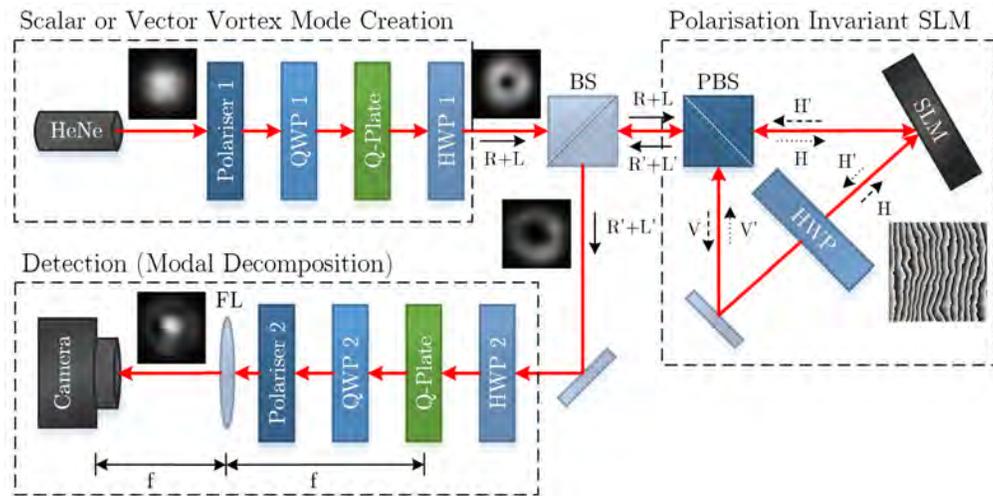


Fig. 2. Simplified schematic diagram of the experimental setup showing the three main parts. The configuration of the half- and quarter-wave plates (HWP and QWP) and polarisers for different modes is summarised in the Appendix. The polarisation invariant SLM is made up by a polarising beam splitter (PBS), mirror and half wave plate rotated to 45° so that arbitrarily polarised scalar and vector beams can be modulated by the SLM, which is encoded with random Kolmogorov turbulence screens. The perturbed beams which return along the same path as the incoming beams are directed by a beam splitter (BS) to the detection part of the setup which performs modal decomposition.

by inverting the creation stage (exploiting the reciprocity of light) followed by an inner product measurement to quantitatively infer the mode coupling weightings by using a technique known as modal decomposition [4, 36].

Each input mode was perturbed by one hundred discrete instances of thin phase atmospheric turbulence of a specific strength. The strength of turbulence was increased in increments of 0.1 from $SR = 0.1$ to 1.0. The channel matrices, M_{scalar} and M_{vector} , described in Section 2 were then generated for each turbulence strength using the averaged and normalised intensity data for comparison to the theoretical calculations in Eq. (8) and (9).

4. Results and discussion

The experimental setup was verified by encoding the SLM with $SR = 1.0$ turbulence, which is simply a grating. This is the control measurement for the experiment and the crosstalk was zero, as expected, shown in Fig. 3.

Figure 3 also shows the channel matrix comparison for $SR = 0.6$ and 0.2 of scalar and vector cases. These two Strehl Ratio's are arbitrary but were chosen because they demonstrate cases for medium and strong turbulence and the clear effect of stronger turbulence, where power is spread from the signal on the diagonal into the crosstalk elements.

The distribution of crosstalk visible in Fig. 3 is very similar to the theory presented in Eqs. (8) and (9). However, according to the theory and the assumption that crosstalk does not occur across polarisation states, some elements in the matrices shown in Fig. 3 should be dark blue (zero) when in fact they show a small amount of signal. This is attributed to the normalisation of the low dynamic range signal in strong turbulence and the efficiency of the PBS.

In order to compare the scalar and vector cases against each other, the crosstalk elements for each input mode were added to each other according to Eq. (13), similar to the theory in

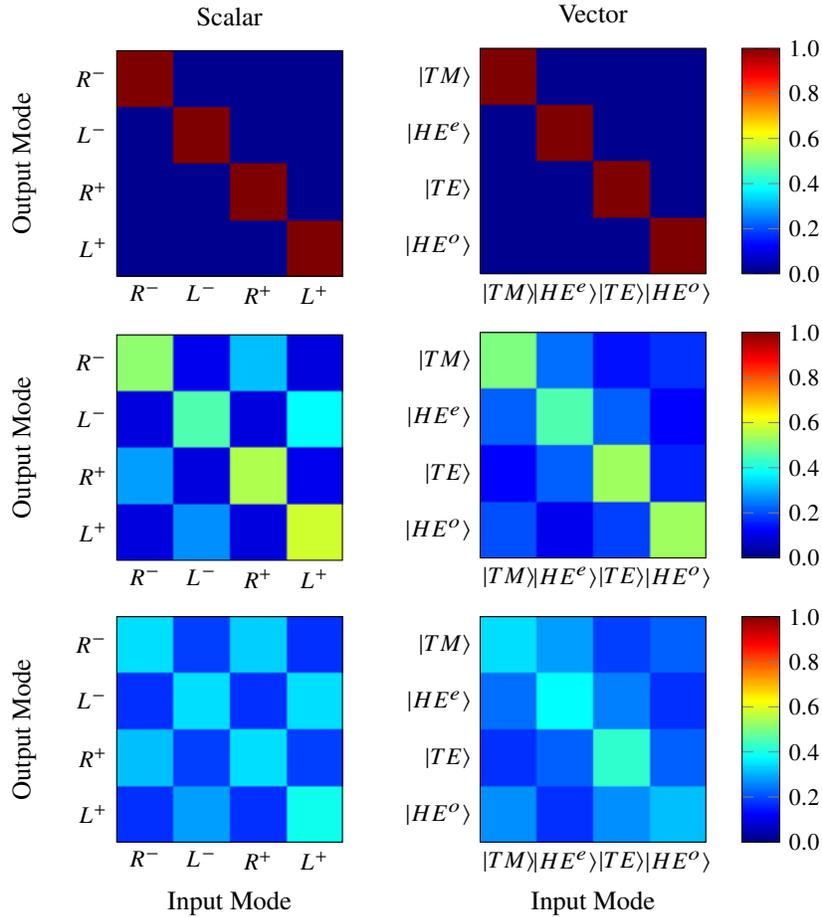


Fig. 3. Comparison of vector and scalar crosstalk for $SR = 1.0$ (top) for experimental setup validation and $SR = 0.6$ (middle) and $SR = 0.2$ (bottom).

Section 2.1. Since each row of the crosstalk matrix is already normalised, the sum must be re-normalised to a percentage, resulting in a total average crosstalk percentage for each turbulence strength for each mode set, visible in Fig. 4.

$$C = \frac{1}{4} \times \sum_{i \neq j} M \times 100\%, \quad (13)$$

where M is the measured scalar or vector channel matrix and C is the associated crosstalk percentage. The crosstalk performance for both scalar and vector modes is identical within the measurement error for all turbulence strengths. These results agree well with the theory presented in Section 2 and indicate that there is no performance benefit to using CVV modes over scalar vortex modes in a thin phase turbulence regime with no birefringence. It should be noted that a similar approach may be used for other basis sets of vector and scalar modes, however, the results in this paper cannot be extended to modes with asymmetric spatial indices such as $|\ell\rangle |R\rangle + |p\rangle |L\rangle$, for example, as their resilience to turbulence cannot be assumed to be equal. Additionally, while it is expected that scalar and CVV modes with higher topological charge will continue to have identical overall crosstalk performance as the theory is general, the absolute values of the crosstalk percentage will vary [37].

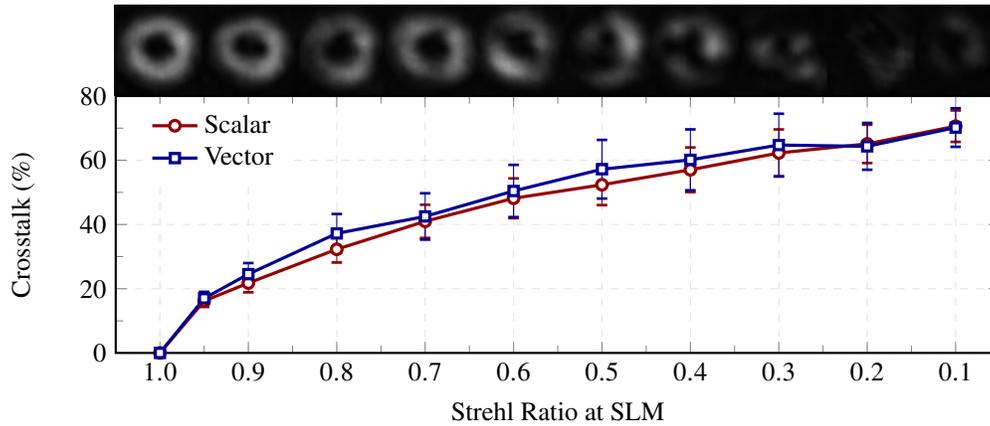


Fig. 4. Percentage of the signal in the crosstalk of scalar and vector cases showing insets of the beam with increasing turbulence. Both mode sets have the same crosstalk performance. Error bars are the standard error of the mean.

Although the overall crosstalk performance for scalar and CVV beams is identical, the crosstalk distribution across modes within the bases are different. It is clearly visible in Fig. 3 that the polarisation component of the scalar basis strongly limits crosstalk between certain modes, however in the vector basis this is not the case. Consequently, the use of CVV modes may in fact limit the possible information capacity of a communication system.

5. Conclusion

It has been inferred in the literature that vector vortex modes are more resilient to atmospheric turbulence than their scalar counterparts. Here we define two similar scalar and cylindrical vector vortex mode bases and theoretically show that their crosstalk in turbulence is the same. This result was then verified experimentally in the thin phase, Kolmogorov regime with various turbulence strengths. The experimental results show identical crosstalk performance within the experimental error.

Appendix

A. Generation of scalar and CVV modes

The generation of both scalar and CVV modes was accomplished using a q -plate ($q = 1/2$) in combination with quarter- and half- wave plates according to the unitary operation \hat{Q} of a q -plate:

$$\hat{Q} \begin{bmatrix} |R\rangle \\ |L\rangle \end{bmatrix} = \begin{bmatrix} |+2q\rangle |L\rangle \\ |-2q\rangle |R\rangle \end{bmatrix}, \quad (14a)$$

$$\hat{Q} \begin{bmatrix} |H\rangle \\ |V\rangle \end{bmatrix} = \begin{bmatrix} |+2q\rangle |R\rangle + |-2q\rangle |L\rangle \\ |+2q\rangle |R\rangle - |-2q\rangle |L\rangle \end{bmatrix}, \quad (14b)$$

where q is the charge of the topological singularity of the q -plate, related to the OAM by $\ell = 2q$. When the input beam is circularly polarized, the output is a scalar beam of the opposite polarization embedded with an OAM of $\pm\ell$. Alternatively, when the input is linearly polarised (horizontal or vertical) the output is a CVV, in which, polarization and OAM are coupled. Experimentally, both sets of modes were generated using a q -plate, linear polarisers, half- and

quarter-wave plates orientated at angles as summarized in Tab. 1. The detection of the modes was achieved by reversing the generation process using a similar configuration and similar optical elements, the orientation angles of the same are also summarised in this table.

Table 1. Configuration of the various components of the experimental setup, depending on the required mode to be generated or decomposed.

Mode	Polariser 1 & 2	HWP 1 & 2	QWP 1 & 2
$ 1\rangle R\rangle$	H	0°	45°
$ 1\rangle L\rangle$	H	-	45°
$ -1\rangle R\rangle$	V	-	45°
$ -1\rangle L\rangle$	H	0°	45°
$ TM\rangle$	H	-	-
$ TE\rangle$	V	-	-
$ HE^e\rangle$	H	0°	-
$ HE^o\rangle$	V	0°	-

Funding

National Research Foundation (NRF); the Claude Leon Foundation; Royal Academy of Engineering; CONACyT; EPSRC Global Challenges Fund (EP/P510968/1).

Acknowledgments

We would like to thank Bienvenu Ndagano for many fruitful conversations as well as Lorenzo Marrucci for providing the q -plates.